

Convergence of the Epstein-Glaser S-matrix in the Sine-Gordon model

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¹Joint work with Dorothea Bahns.

Outline of the talk

- 1 Algebraic QFT and its generalizations
 - Outline of the pAQFT framework
 - Scalar field

- 2 The Sine-Gordon model

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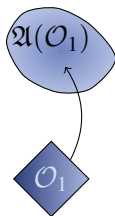
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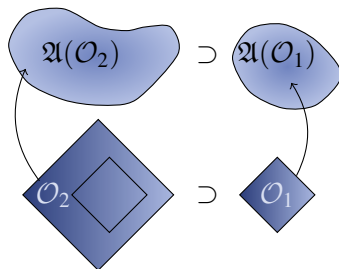
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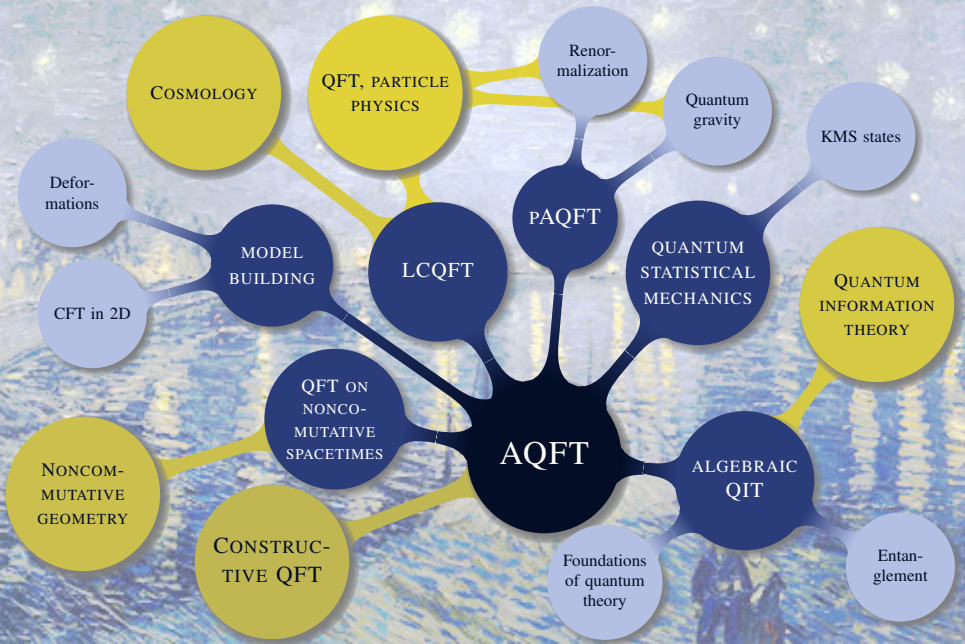
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- The physical notion of subsystems is realized by the condition of **isotony**, i.e.: $\mathcal{O}_1 \subset \mathcal{O}_2 \Rightarrow \mathfrak{A}(\mathcal{O}_1) \subset \mathfrak{A}(\mathcal{O}_2)$. We obtain a **net of algebras**.



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- Mathematically, AQFT makes use of **functional analysis** techniques (operator algebras), but its various generalizations involve many other branches of mathematics.

Different aspects of AQFT and relations to physics



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 - Generalization to gauge theories using homological algebra ([Hollands 08, Fredenhagen-KR 11]).

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- **Dynamics**: we use a modification of the Lagrangian formalism. Since the manifold M is non-compact, we need to introduce a cutoff function into the action functional. For the free scalar field

$$S_M(f)(\varphi) = \frac{1}{2} \int (\nabla_\mu \varphi \nabla^\mu \varphi - m^2 \varphi^2)(x) f(x) d\mu(x).$$

In general an action is a map $S_M : \mathcal{D}(M) \rightarrow \mathcal{C}^\infty(\mathcal{E}(M), \mathbb{R})$, where $\mathcal{D}(M) \equiv \mathcal{C}_c^\infty(M, \mathbb{R})$ are compactly supported smooth functions.

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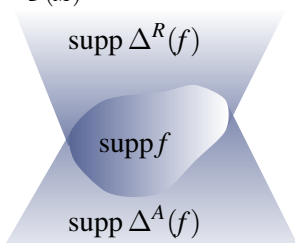
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- Under some technical assumptions on M , P admits retarded and advanced Green's functions Δ^R, Δ^A . They satisfy:

$$P \circ \Delta^{R/A} = \text{id}_{\mathcal{D}(M)}, \Delta^{R/A} \circ (P|_{\mathcal{D}(M)}) = \text{id}_{\mathcal{D}(M)}$$
and

$$\text{supp}(\Delta^R) \subset \{(x, y) \in M^2 | y \in J_-(x)\},$$

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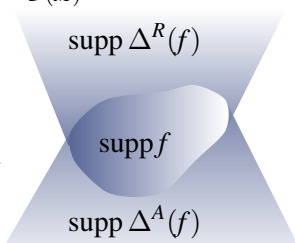
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$$\begin{aligned} \text{supp}(\Delta^R) &\subset \{(x, y) \in M^2 \mid y \in J_-(x)\}, \\ \text{supp}(\Delta^A) &\subset \{(x, y) \in M^2 \mid y \in J_+(x)\}. \end{aligned}$$

- Their difference is the causal propagator

$$\Delta \doteq \Delta^R - \Delta^A.$$



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- The free QFT is defined as $\mathfrak{A}_0(M) \doteq (\mathcal{F}(M)[[\hbar]], \star, *)$, where $F^*(\varphi) \doteq \overline{F(\varphi)}$ and $\mathcal{F}(M)$ is an appropriate functional space (some WF set conditions on $F^{(n)}(\varphi)$ s induced by W).

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- Define the time-ordered product $\cdot_{\mathcal{T}}$ on $\mathcal{F}_{\text{reg}}(M)[[\hbar]]$ by:

$$F \cdot_{\mathcal{T}} G \doteq \mathcal{T}(\mathcal{T}^{-1}F \cdot \mathcal{T}^{-1}G)$$

Interaction

- We now have an algebraic structure with two products $(\mathcal{F}_{\text{reg}}(M)[[\hbar]], \star, \cdot_{\mathcal{T}})$, where \star is non-commutative, $\cdot_{\mathcal{T}}$ is commutative and they are related by a causal relation:

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- We define the **interacting** star product as:

$$F \star_{\text{int}} G \doteq R_V^{-1} (R_V(F) \star R_V(G)),$$

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- To extend \mathcal{T}_n to arbitrary local functionals we use e.g. the causal approach of Epstein and Glaser.

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- The 2-point function of the free massless scalar field in 2D coincides with the Hadamard parametrix

$$W(x) = -\frac{1}{4\pi} \ln \left(\frac{-x \cdot x + i\epsilon t}{\Lambda^2} \right)$$

where $\Lambda > 0$ is the scale parameter.

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- In particular, for $\mathbf{x} = \mathbf{y}$ and $t > t'$,

$$V_a(t, \mathbf{x}) \star V_{a'}(t', \mathbf{x}) = e^{aa' i\hbar/2} V_{a'}(t', \mathbf{x}) \star V_a(t, \mathbf{x}),$$

which is the well-known braiding property for vertex operators.

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- Hence \star and \star_v are equivalent products, and α_v is a “gauge transformation”.

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- Similar for the S -matrix:

$$\omega_v(\mathcal{S}(\lambda :V:_v)) \doteq \alpha_v \left(e_{\mathcal{T}}^{i\lambda :V:_v / \hbar} \right) (0) = e_{\mathcal{T}_v}^{i\lambda V / \hbar}(0).$$

Here $\cdot_{\mathcal{T}_v}$ is the time-ordered product corresponding to \star_v .

Convergence of the S-matrix

Theorem (Bahns, KR 2016)

The formal S-matrix $\alpha_v \circ \mathcal{S}(\lambda : V : v) = e^{\frac{i\lambda V}{\hbar}}$ in the Sine-Gordon model with $V = \frac{1}{2}(V_a(f) + V_{-a}(f))$ and $0 < \beta = \hbar a^2 / 4\pi < 1$, $f \in \mathcal{D}(M)$, converges as a functional on the configuration space in the appropriate topology (related to Hörmander topology on distribution spaces).

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- Direct proof of the convergence of the S-matrix in the Minkowski signature.

Convergence of the S-matrix

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The formal S-matrix $\alpha_v \circ \mathcal{S}(\lambda:V:v) = e_{\mathcal{T}_v}^{i\lambda V/\hbar}$ in the Sine-Gordon model with $V = \frac{1}{2}(V_a(f) + V_{-a}(f))$ and $0 < \beta = \hbar a^2/4\pi < 1$, $f \in \mathcal{D}(M)$, converges as a functional on the configuration space in the appropriate topology (related to Hörmander topology on distribution spaces).

- Direct proof of the convergence of the S-matrix in the Minkowski signature.
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- Direct proof of the convergence of the S -matrix in the Minkowski signature.
- No issues with positivity/IR problems, no Wick rotation.
- The abstract formal S -matrix is constructed before a state is chosen.

Some details of the proof I

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- Hence the estimates boil down to estimates of expressions of the form:

$$\prod_{1 \leq i < j \leq k} |\tau_{ij}^2 - \zeta_{ij}^2|^\beta \prod_{1 \leq i \leq k, k < j \leq n} |\tau_{ij}^2 - \zeta_{ij}^2|^{-\beta} \prod_{k < i < j \leq n} |\tau_{ij}^2 - \zeta_{ij}^2|^\beta.$$

with the time variable differences $\tau_{ij} = t_i - t_j$ and the space variable differences $\zeta_{ij} = \mathbf{x}_i - \mathbf{x}_j$.

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- After a change of variables, this expression can be rewritten as a determinant (analogous to [Fröhlich 76]):

$$D_{ij} = \begin{cases} w_j^{i-1} & , \quad 1 \leq i \leq l - k , \\ 1/(z_{i-l+k} - w_j) & , \quad l - k < i \leq l . \end{cases}$$

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- The estimates are possible due to the fact that $\beta < 1$ and for the Vandermonde determinants they require one to choose the support of the cutoff function sufficiently small.
- The later requirement would not be necessary if we were using a singular state obtained as the limit of the massive theory, instead of a Hadamard state.
- In our future work we expect to be able to drop the condition on the support of the test function, in an appropriate class of Hadamard states.

Constructing a net of von Neumann algebras

- Use the Bogoliubov formula to construct the interacting fields:

$$\begin{aligned} R_{\lambda V}(F) &= -i\hbar \frac{d}{d\mu} \mathcal{S}(\lambda : V :)^{-1} \mathcal{S}(\lambda : V : + \mu : F :) \Big|_{\mu=0} \\ &\equiv -i\hbar \frac{d}{d\mu} \mathcal{S}_{\lambda : V :}(\mu : F :) \Big|_{\mu=0}, \end{aligned}$$

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- Prove covariance.
- Construct the local net using the prescription given in [Fredenhagen, KR 2015].

Outlook

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- Extend our results to some range of $\beta > 1$ (e.g. super-renormalizable range).
- Construct conserved currents for Sine Gordon (i.e. show integrability).
- Apply the same methods to a larger class of integrable models.
- Show equivalence with the $\mathcal{O}(3)$ model and the Thirring model.

Last, but not least. . .

For Bernard and Henning:



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For Dorothea, Chris and Gandalf: **Thank you for this wonderful event!**





Thank you very much for your attention!